

$N=2$ SQM

It is based on so-called extended superspace

$t, \theta_+, \theta_-, \bar{\theta}_+, \bar{\theta}_-$

$Q_{\pm} = \frac{\partial}{\partial \theta_{\pm}} + \bar{\theta}_{\pm} \frac{\partial}{\partial t}, \quad \bar{Q}_{\pm} = \frac{\partial}{\partial \bar{\theta}_{\pm}} + \theta_{\pm} \frac{\partial}{\partial t}$

$D_{\pm} = \frac{\partial}{\partial \theta_{\pm}} - \bar{\theta}_{\pm} \frac{\partial}{\partial t}, \quad \bar{D}_{\pm} = \frac{\partial}{\partial \bar{\theta}_{\pm}} - \theta_{\pm} \frac{\partial}{\partial t}$

$\{Q_+, \bar{Q}_+\} = \{Q_-, \bar{Q}_-\} = \frac{\partial}{\partial t} \leftarrow$ to be 4 SUSY of the theory

$\{D_+, \bar{D}_+\} = \{D_-, \bar{D}_-\} = \frac{\partial}{\partial t}$

other anticommutators of Q 's and D 's are equal to zero.

New concept - chiral superfields.

$X'(t, \theta_{\pm}, \bar{\theta}_{\pm})$ is called chiral if it is annihilated by \bar{D}_{\pm} . $\bar{D}_{\pm} X' = 0$ example

$\left(\frac{\partial}{\partial \bar{\theta}_+} - \theta_+ \frac{\partial}{\partial t} \right) f(t, \theta_+, \bar{\theta}_+) = 0 \quad (\epsilon)$

θ_+ is \bar{D} chiral

t is not chiral, $\bar{\theta}_+$ is not chiral

$t_c = t + \bar{\theta}_+ \theta_+$ is chiral

we may solve equation (ϵ)

considering $F(\theta_+, t_c) \rightarrow$ is

\bar{D} -chiral.

The chiral field is a function of θ_+, θ_-

and $t_c = t + \bar{\theta}_+ \theta_+ + \bar{\theta}_- \theta_-$

Solution to chirality equations

$\hat{\Phi}^i = X^i(t_c) + \psi_+^i(t_c)\theta_+ + \psi_-^i(t_c)\theta_- + F^i\theta_+\theta_-$
 How Q and \bar{Q} are acting on chiral fields. Action of Q_{\pm} is easy to compute - really, \bar{Q}_{\pm} is different from

\bar{D}_{\pm} by a sign

$$Q_+ = \frac{\partial}{\partial \theta_+} + \bar{\theta}_+ \frac{\partial}{\partial t_c}$$

$$Q_{\pm} F(t_c, \theta_+, \theta_-) = \frac{\partial F}{\partial \theta_{\pm}} +$$

$$\bullet \bar{\theta}_+ \frac{\partial F}{\partial t_c} + \bar{\theta}_+ \frac{\partial F}{\partial t_c} = \frac{\partial F}{\partial \theta_{\pm}}$$

$$\bar{Q}_{\pm} = -\bar{D}_{\pm} + 2\theta_{\pm} \frac{\partial}{\partial t}$$

$$\bar{Q}_{\pm} F = 2\theta_{\pm} \frac{\partial F}{\partial t}, \text{ so there is}$$

the following representation

$$Q_{\pm} \rightarrow \frac{\partial}{\partial \theta_{\pm}} \quad \bar{Q}_{\pm} = 2\theta_{\pm} \frac{\partial}{\partial t_c}$$

they actually anticommute as they should

Action of Q_{\pm} on components:

$Q_+ \hat{\Phi}^i = \psi_+^i + \theta_- F^i$	$\delta_{Q_+} X^i = \psi_+^i$	$\delta_{Q_+} \psi_-^i = F^i$
$Q_- \hat{\Phi}^i = \psi_-^i - \theta_+ F^i$	$\delta_{Q_-} X^i = \psi_-^i$	$\delta_{Q_-} \psi_+^i = -F^i$

$$\bar{Q}_+ \hat{\Phi}^i = \theta_+ \frac{\partial X^i}{\partial t} + \theta_+ \theta_- \frac{\partial \psi_-^i}{\partial t} \quad \left| \quad \delta_{\bar{Q}_+} \psi_+^i = \frac{\partial X^i}{\partial t} \right.$$

$$\delta_{\bar{Q}_+} F^i = \frac{\partial F^i}{\partial t}$$

similarly $\delta_{\bar{Q}_-} \psi_-^i = \frac{\partial X^i}{\partial t}$

$$\delta \bar{Q}_- F^i = \frac{\partial \Psi^i}{\partial t}$$

Let us check the pair: δQ_+ , $\delta \bar{Q}_+$:

$$\left\{ \frac{\delta}{\delta Q_+}, \frac{\delta}{\delta \bar{Q}_+} \right\} X^i = \frac{\partial X^i}{\partial t}$$

$N=2$ SQM we also used antichiral fields

$$\bar{\Phi}^i: D_+ \bar{\Phi}^i = 0$$

$$\bar{\Phi} = f(t_a, \bar{\Theta}_+, \bar{\Theta}_-)$$

$$t_a = t + \theta_+ \bar{\Theta}_+ + \theta_- \bar{\Theta}_-$$

$$\begin{aligned} \bar{\Phi}^i &= \bar{X}^i(t_a) + \bar{\Psi}_+^i(t_a) \bar{\Theta}_+ + \bar{\Psi}_-^i(t_a) \bar{\Theta}_- + \\ &\quad \bar{F}(t_a) \bar{\Theta}_+ \bar{\Theta}_- \end{aligned}$$

The role of Q_+ and \bar{Q}_+ would be interchanged: exchange $\theta \leftrightarrow \bar{\theta}$, $Q \leftrightarrow \bar{Q}$.

Action - invariant of the supertranslation

in $N=1$ case we had $\int dt \int d\theta d\bar{\theta} W(\Phi)$
 in $N=2$ case we have similarly the following terms:

$$\begin{aligned} &\int dt_c \int d\theta_+ d\bar{\theta}_- W(\Phi_c) = \\ &\int dt \left(\frac{\partial W}{\partial \Phi^i} F^i + \frac{\partial^2 W}{\partial \Phi^i \partial \Phi^j} \Psi_+^i \Psi_-^j \right) \end{aligned}$$

if $W = \Phi$ it is just $\int dt F^+$ and

is called F-term

$$\text{similarly } \int dt_a \int d\bar{\theta}_+ d\bar{\theta}_- W(\bar{\Phi}_a) = \\ = \int dt \left(\frac{\partial W}{\partial \bar{\Phi}^i} \bar{F}^i + \frac{\partial^2 W}{\partial \bar{\Phi}^i \partial \bar{\Phi}^j} \bar{\Psi}_+^i \bar{\Psi}_-^j \right)$$

Before, in $N=1$ SQM we got kinetic term from $g_{ij} D\Phi^i D\Phi^j \leftrightarrow$ but not now!

$$\int dt d\theta_+ d\theta_- d\bar{\theta}_+ d\bar{\theta}_- K(\Phi_c, \bar{\Phi}_a) = ?$$

simplest case $K = \Phi_c \bar{\Phi}_a$, I need 4- θ 's
easiest structure is - get 2θ from $F\theta_+ \theta_-$ on Φ_c
and $2\bar{\theta}$ from $\bar{\theta}_+ \bar{\theta}_-$ on $\bar{\Phi}_a$: $\int F \bar{F} dt$
(it is an analogue of $N=1$ SQM).
 F^2 term in

Now we need derivatives

Take $\theta_+ \psi_+(t_c)$ and $\bar{\theta}_+ \psi_+(t_a)$
from Φ_c from $\bar{\Phi}_a \rightarrow$

we already have $\theta_+ \bar{\theta}_+$, how to get $\bar{\theta}_- \theta_-$?
 $\bar{\theta}_- \theta_-$? from the difference in $\theta_- \bar{\theta}_-$
dependence of t_c and t_a !

I may do tedious calculations exp.

$$\psi_+(t_c) \text{ as } \psi_+(t) + \bar{\theta}_- \theta_- \frac{\partial \psi_+}{\partial t} + \dots \\ \bar{\psi}_+(t_a) \text{ as } \bar{\psi}_+(t) + \theta_- \bar{\theta}_- \frac{\partial \bar{\psi}_+}{\partial t} + \dots$$

$\psi_+ \frac{\partial \bar{\psi}_+}{\partial t} + \bar{\psi}_+ \frac{\partial \psi_+}{\partial t}$ - not a total derivative due to

fact that $\psi_+, \bar{\psi}_+$ are anticommuting
kinetic term for fermions, sim for

$\psi_-, \bar{\psi}_-$

what about bosons

$X(t_c) \bar{X}(t_a)$ - similarly, \bar{I} expand
to get $\theta \bar{\theta}$ from $\theta \bar{\theta}$ dependence of t_c
and t_a . But now \bar{I} get 2 time
derivatives. $\int dt (\partial_t X \partial_t \bar{X})$

(after we eliminate total derivatives
on t).

So we get the action in components
similar to $N=1$ SQM, with the dif.
that we have COMPLEX STRUCTURE
ON THE TARGET.

Hol. coordinates are chiral fields,
antihol. coordinates are antichiral
fields.

Remark, in $N=1$ theory we
may try to int. only one constr.

$\bar{D} f_c = 0$, then field would go

$f_c = \psi + \theta \tilde{F}$; how to write

kinetic term? $\int d\theta \int W(f_c)$?

probably, it is a way to 1-order formulation of

SQM. - \bar{I} never tried. -
- good question

$$\int d^2\theta_+ d^2\theta_- K(\Phi_a^i, \bar{\Phi}_c^{\bar{j}}) = \frac{\partial^2 K}{\partial\Phi^i \partial\bar{\Phi}^{\bar{j}}} \chi^i \chi^{\bar{j}} + \dots$$

D-term

$$\frac{\partial^2 K}{\partial\Phi^i \partial\bar{\Phi}^{\bar{j}}} \leftrightarrow g_{i\bar{j}}$$

- metric on the target space.

But is exactly the condition for metric to be Kähler.

Really, $\partial_i g_{j\bar{k}} = \partial_{\bar{j}} g_{i\bar{k}} \rightarrow$

$$g_{i\bar{k}} = \partial_i (V_{\bar{k}}) \quad (1)$$

similarly, if we also impose

$$\partial_{\bar{i}} g_{j\bar{k}} = \partial_{\bar{k}} g_{j\bar{i}} \rightarrow$$

$$g_{i\bar{k}} = \partial_{\bar{k}} (\tilde{V}_{\bar{i}}) \quad (2)$$

$$(1) \& (2) \Rightarrow g_{i\bar{j}} = \frac{\partial K}{\partial\chi^i \partial\chi^{\bar{j}}}$$

exactly what we have from the D-term